

Test Beam Line (TBL)

Beam Dynamics studies

E. Adli, CERN AB/ABP / UiO
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The purpose of the TBL

- Validate the drive beam decelerator concept
 - Demonstrate the efficiency of RF power production
 - Demonstrate the stability of the drive beam
 - Demonstrate algorithms and technology for CLIC

Purpose and content of presentation

Purpose of presentation:

“try to outline alignment precision requirements, starting from the energy extraction requirement”

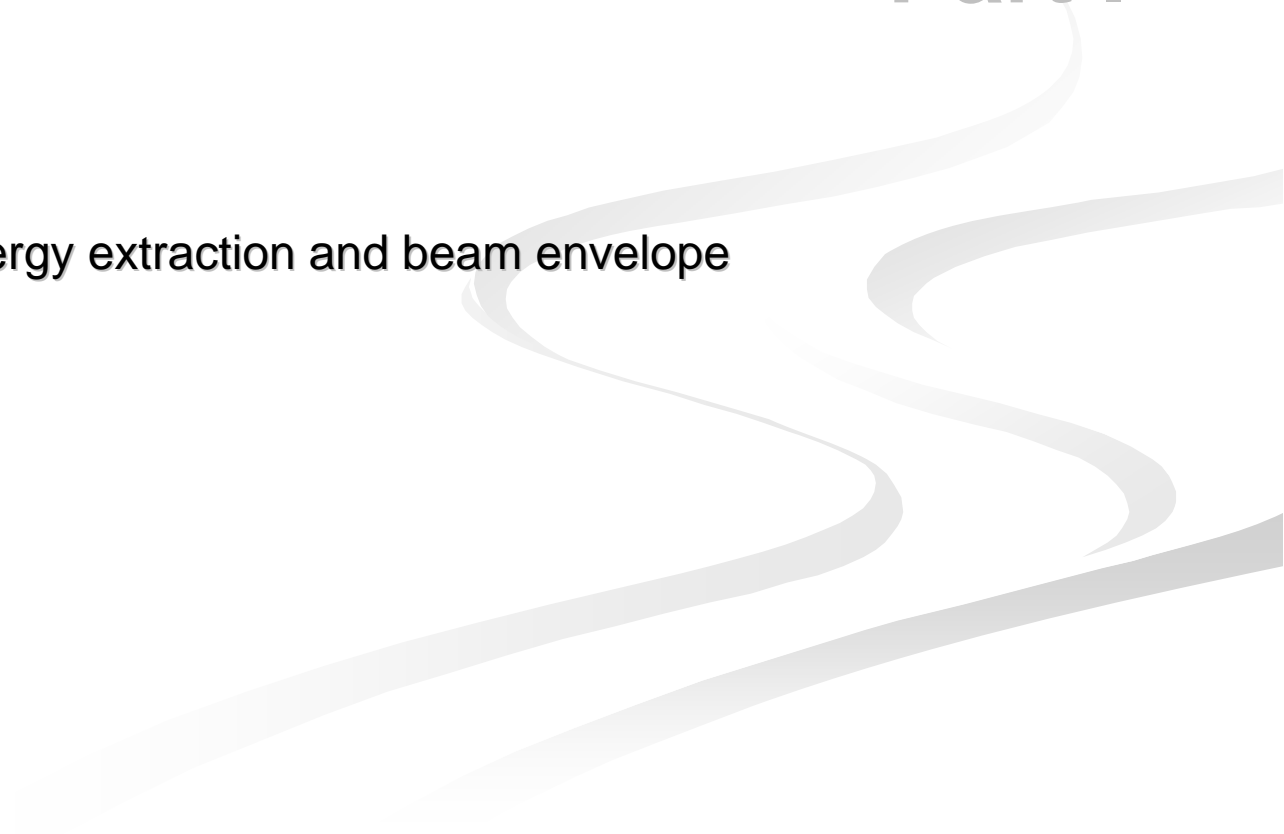
Contents:

1. Energy extraction and beam envelope
2. Beam stability with transverse wakes (summary)
3. Effect of component misalignments

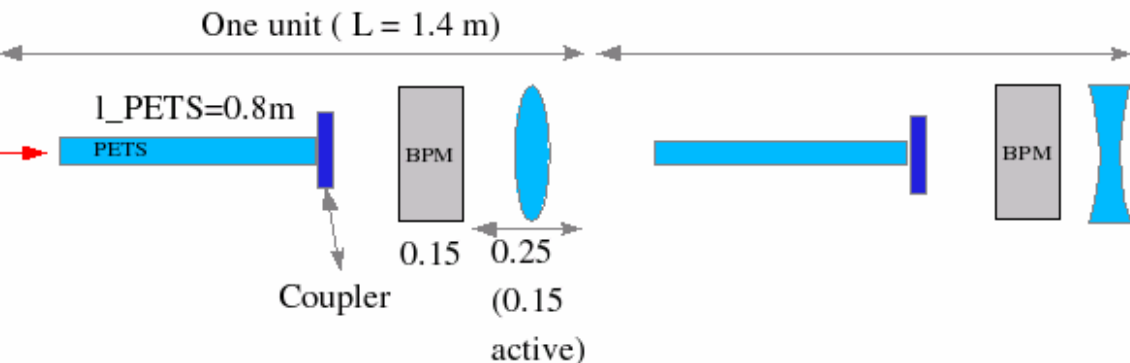
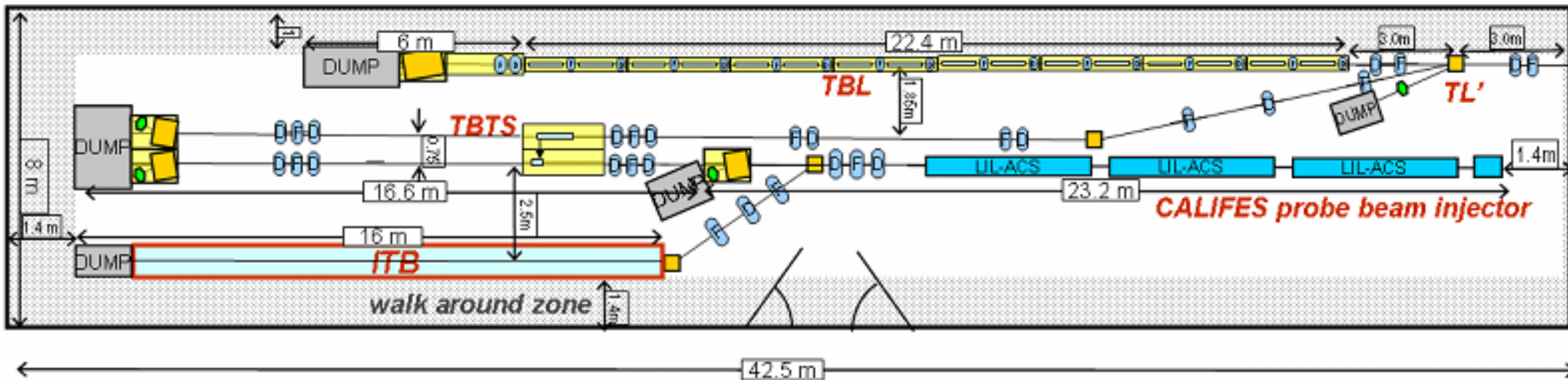
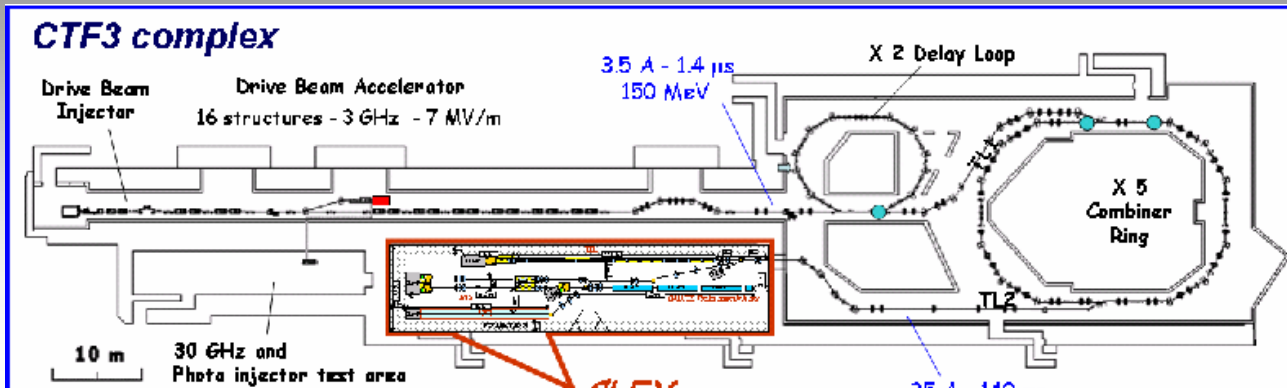
“CLIC” versus “TBL”: the graphs and numbers here are for TBL, unless other stated, but the principles are the same for CLIC (one deceleration station)

Part I

Energy extraction and beam envelope

The background features several thick, light gray wavy lines that flow from the bottom left towards the top right, creating a sense of movement and depth.

Simulation set-up: LATTICE



Simulation lattice:

16 units of one of each:

- PETS (coupler as drift)
- Quad
- BPM

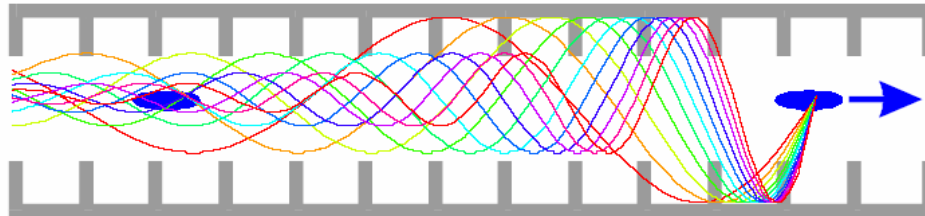
The CTF3 drive beam

- High-current, low-energy beam for strong wake field generation
- Initial beam parameters used for these simulations:
 - $E_0 = 150$ MeV (no energy spread)
 - $I = 30$ A
 - $d = 25$ mm (bunch spacing, $f_b = 12$ GHz)
 - Gaussian bunch, $\sigma_z = 1$ mm
 - $N = 200$ (enough for steady-state situation to be reached).
 - $\varepsilon_N = 150$ μm



Deceleration

- Particles will feel parasitic loss and induce a wake field in the PETS
- The wake field will interact with and further decelerate :
 - 1) rear part of bunch (**single-bunch** effect)
 - 2) following bunches (**multi-bunch** effect)



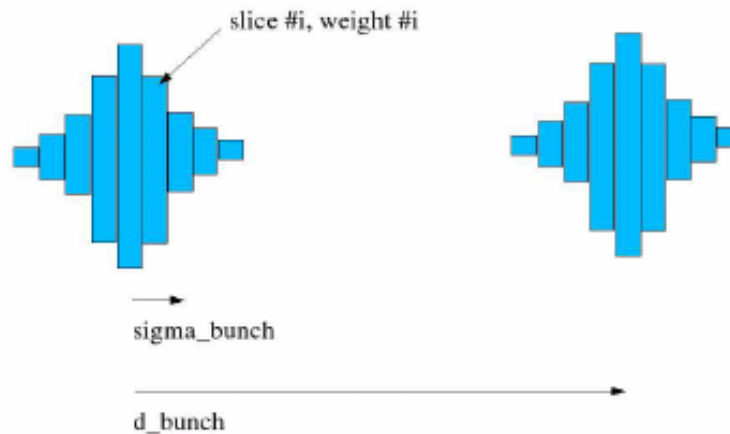
- The integrated effect in a PETS on a witness particle due to a source particle is given by (h.o.m. ignored)

$$\int_0^{l_{cav}} F_L(z) ds \approx -q_s q_w w_L(z)$$

where w_L is the std. longitudinal monopole wake function

Simulation software: PLACET

- The simulation package used here is PLACET (D. Schulte)
- Allows to study the effect of single-bunch + multi-bunch wakes precisely
- Beam model used here: sliced beam with a Gaussian longitudinal profile



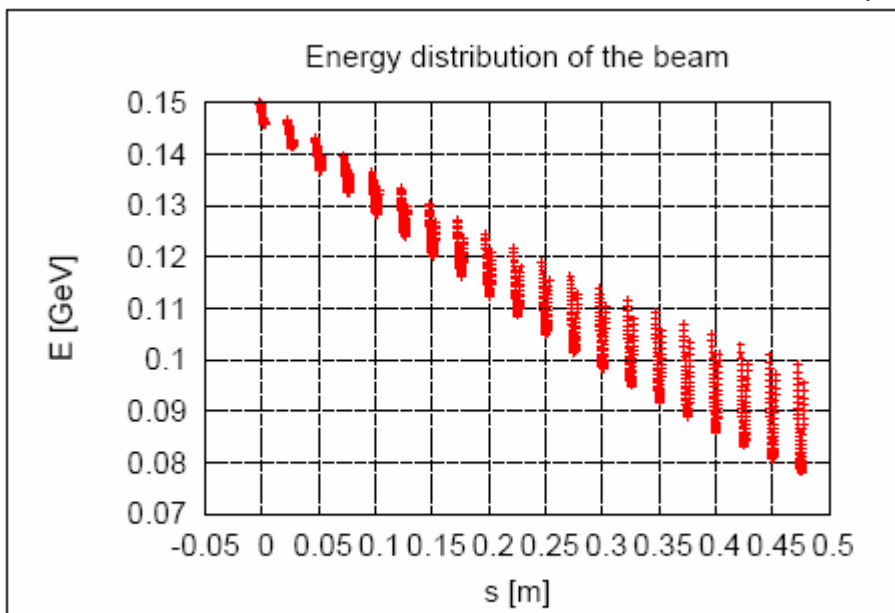
- wake acting on a given slice is simply the sum contribution from all leading slices (multi- and single-bunch effects treated on equal footing)

Simulation results: energy extraction

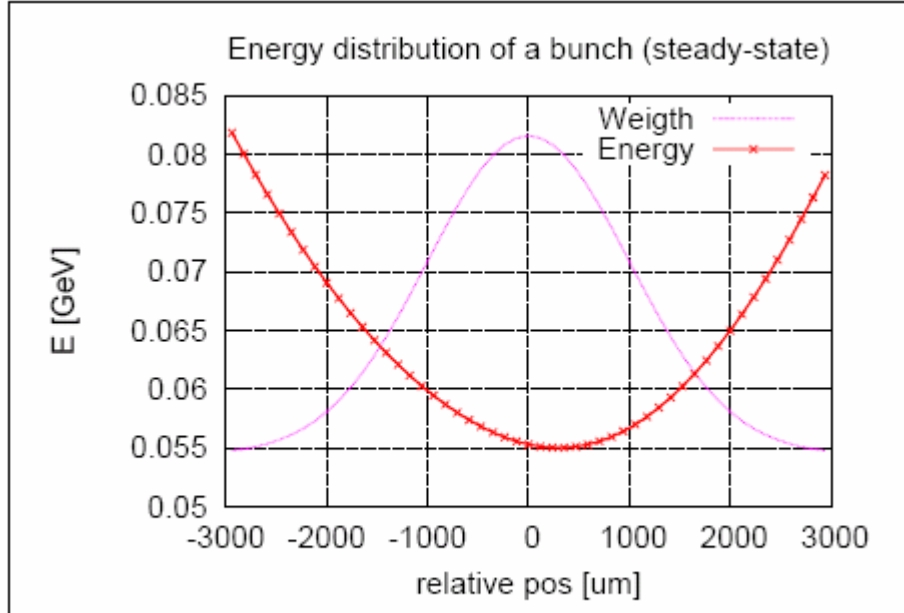
■ PETS longitudinal wake parameters:

- $R'/Q = 2294.7 \text{ } \Omega/\text{m}$ (linac-convention)
- $f_L = 11.99 \text{ GHz}$
- $\beta_g = 0.4529$

■ Beam energy profile after lattice: (initial: flat $E_0 = 150 \text{ MeV}$)



Beam energy after lattice

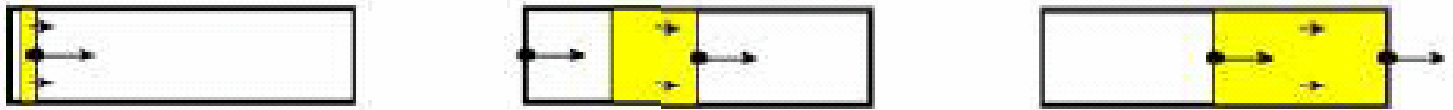


Steady state bunch energy profile

- **NB: start of beam / bunch is to the left! (PLACET output def.)**

Wake calculations and group velocity

- The wake is calculated using GdfidL (I. Syratchev), modeled as a single monopole mode traveling out of the PETS with a high group velocity (β_g) [and extracted to HDS]



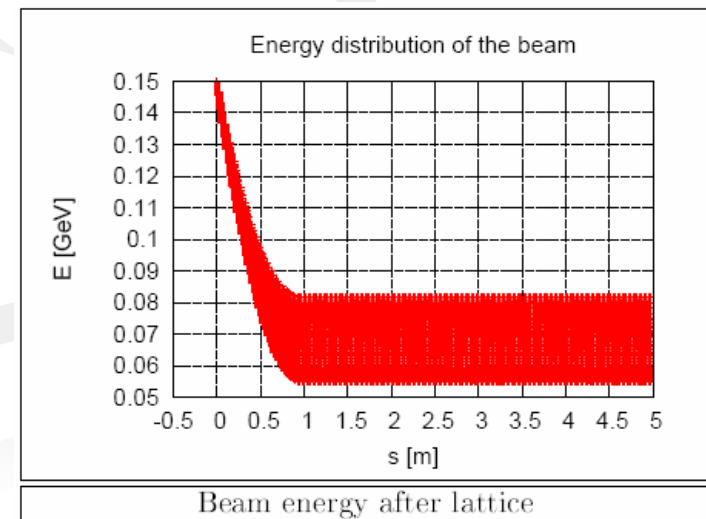
- In the longitudinal wake function this leads to
 - factor $1/(1 - \beta_g)$ (concentration of the field)
 - catch-up distance for the trailing bunch, $s = z\beta_g/(1 - \beta_g)$
- The wake parameters R'/Q , β_g and f are taken as input to PLACET the simulation δ -wake:

$$W_{\delta L}(z) = \omega_L \frac{R'}{Q} \frac{1}{1 - \beta_L} \cos\left(\omega_L \frac{z}{c}\right) \left(L - z \frac{\beta_L}{1 - \beta_L}\right) [V/C]$$

Steady-state

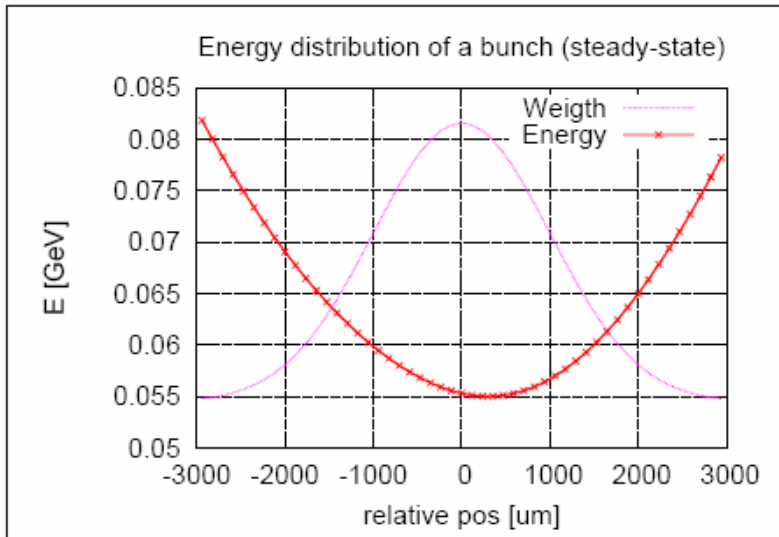
- Catch-up with field from n bunches ahead at a distance $s = nd\beta_g/(1-\beta_g)$
- Steady-state energy profile is thus reached after
$$n = (I_{PETS}/d)(1-\beta_g)/\beta_g = 39 \text{ bunches}$$
- Steady-state power can be calculated as $P = \frac{\omega}{4v_g}(R'/Q)l^2_{PETS}I^2F^2(\sigma)$
($P = 172\text{MW}$, or $P \approx 166\text{MW}$ if wall losses are included)

When the bunch profile and energy extraction efficiency is discussed we always talk about the **steady-state situation**.

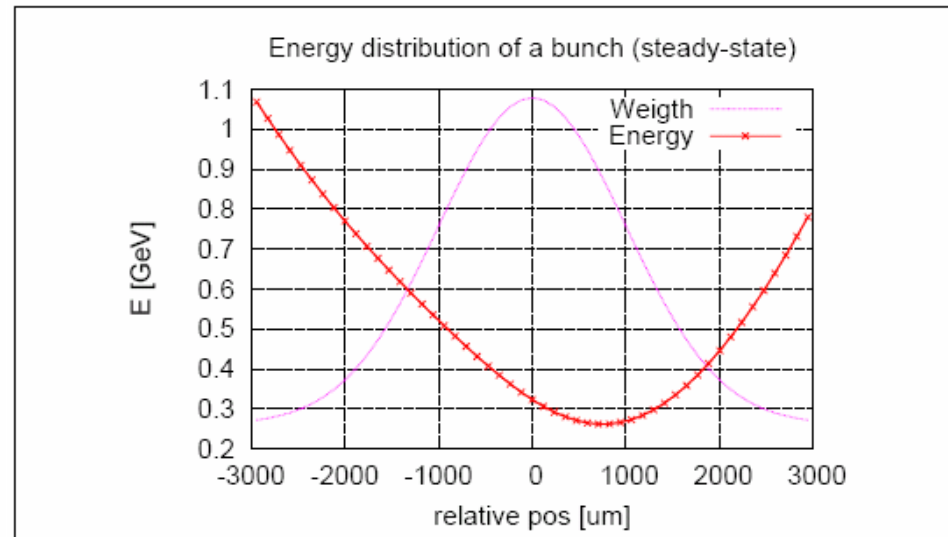


Steady state bunch profile

- The steady-state bunch profile depends on the multi-bunch effects as well as the single-bunch effects
- Multi-bunch wake alone would form a symmetrical energy profile (cosine-like wake function, combined with Gaussian distribution)
- Single-bunch wake: last part of the bunch will be more decelerated than the first -> point of minimum energy shifted towards the end
- However, for our case, $n = (I_{PETS}/d)(1-\beta_g)/\beta_g = 39$ multi-bunch is dominant



Steady state bunch energy profile



CLIC 12 GHz, $\eta_{dist} = 97.3\%$ ($I_{PETS,CLIC} = 0.3I_{PETS,TBL}$)

- Compare with e.g. profile for CLIC 12 GHz ($I_{PETS} = 0.23$ m)

Energy extraction efficiency: η

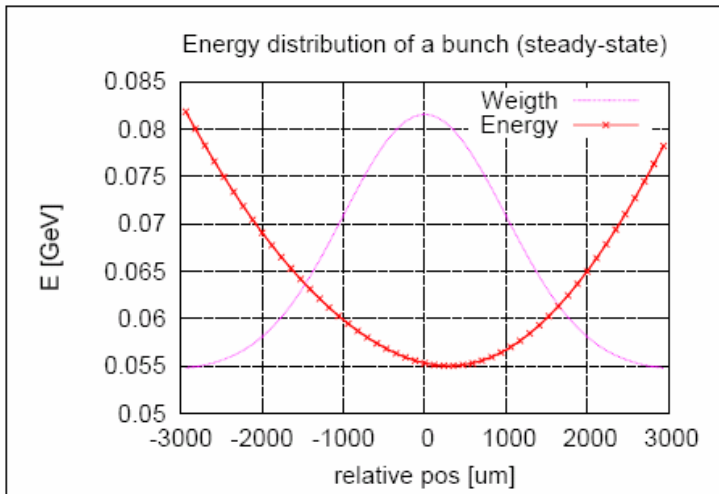
- $\eta = P_{in}/P_{out}$: steady state power extraction eff: $\eta = P[W] \times N / E0[eV] \times I[A]$

- Suggestion: it could be useful to express the extraction efficiency as:

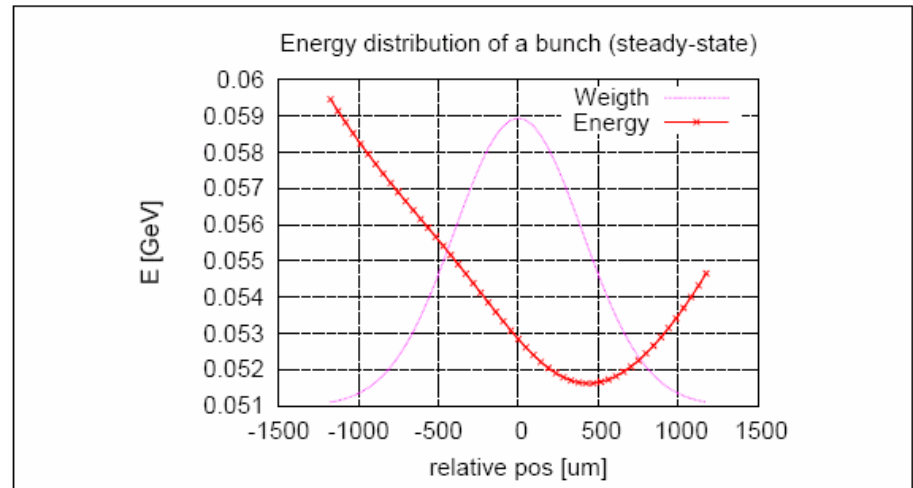
$$\eta = S \times F(\sigma) \times \eta_{dist}$$

where for TBL nominal parameters we get:

- $S = 63.3 \%$ (max energy spread)
- $\eta = S \times F(\sigma) \times \eta_{dist} = 63.3 \% \times 96.9 \% \times 99.9 \% = 61.3 \%$



Steady state bunch energy profile



$\sigma_z = 400 \mu m$ ($\eta = S \times F(\sigma) \times \eta_{dist} = 65.6\% \times 99.5\% \times 99.0\%$)

- (can be changed with detuning: not discussed further here)

The CTF3 drive beam

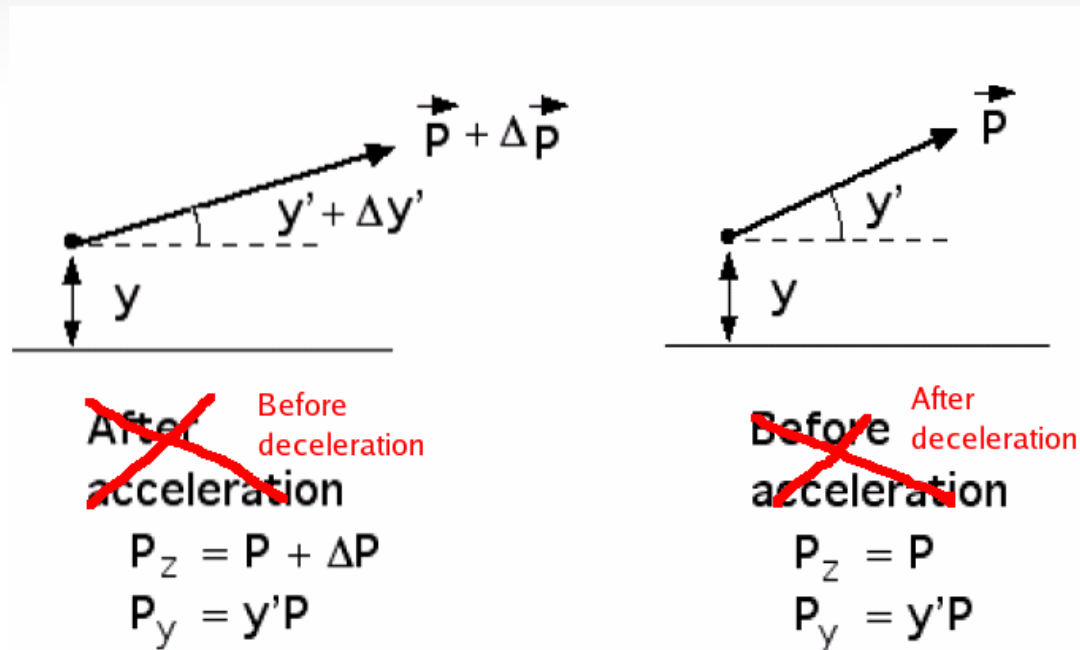
- High-current, low energy beam for strong wake field generation
- Initial beam parameters used for these simulations:
 - $E_0 = 150$ MeV (no energy spread)
 - $I = 30$ A
 - $d = 25$ mm ($f_b = 12$ GHz)
 - Gaussian bunch: $\sigma_z = 1$ mm
 - $N = 200$ (enough for steady-state situation to be reached).
 - $\varepsilon_N = 150$ μm



- **Resulting parameters:**
 - **$P = 166$ MW (steady-state power production)**
 - **$S = 63.3$ % (max .energy spread)**
 - **$\eta = 61.3$ % (steady-state extraction efficiency)**

Energy spread and beam envelope

- Why is the max. energy spread, S , important?
- In the TBL we will have the effect of *adiabatic undamping*



(fig: A. Chao)

- The divergence, $y' = dy/ds$, and thus also the beam envelope will increase with decreasing energy

Calculation of the max. beam envelope

- This implies that as the beam is decelerated its transverse size will grow, even without considering transverse wake kicks or machine imperfections

- The rms beam size is

$$\sigma_{x,y} = \sqrt{\beta_{x,y} \varepsilon_{x,y}} = \sqrt{\beta_{x,y} \varepsilon_{N,x,y} / \gamma}$$

- We define the adiabatic “3-sigma beam envelope” as

$$r_{ad} = \sqrt{3^2 \sigma_x^2 + 3^2 \sigma_y^2}$$

where γ is for the **lowest energy particle** in the bunch

- Setting in for S, with γ_0 the initial gamma we get the value in the middle of a quad:

$$r_{ad} = \sqrt{3^2 (\tilde{\beta} + \hat{\beta}) \varepsilon_N / (1-S) \gamma_0} \approx 3 \cdot 2 \sqrt{L_{unit} \varepsilon_N / (1-S) \gamma_0}$$

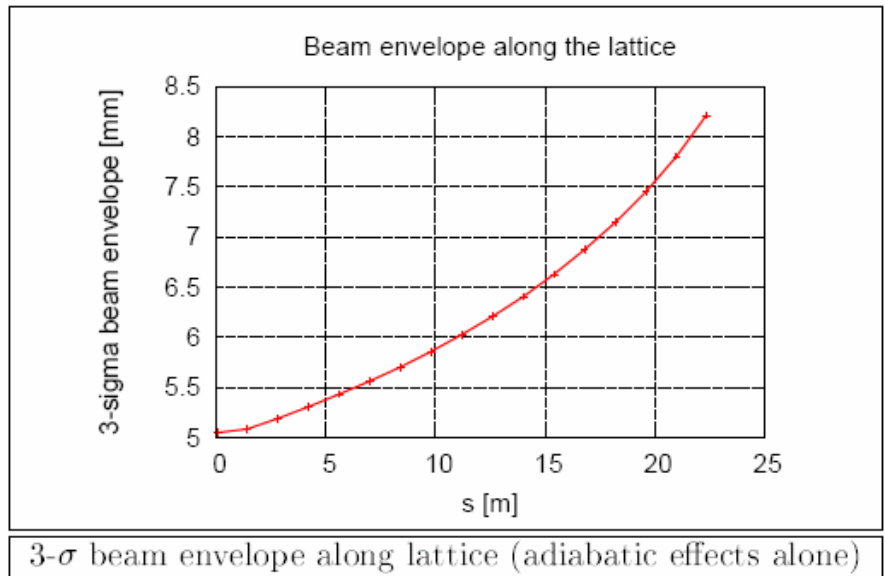
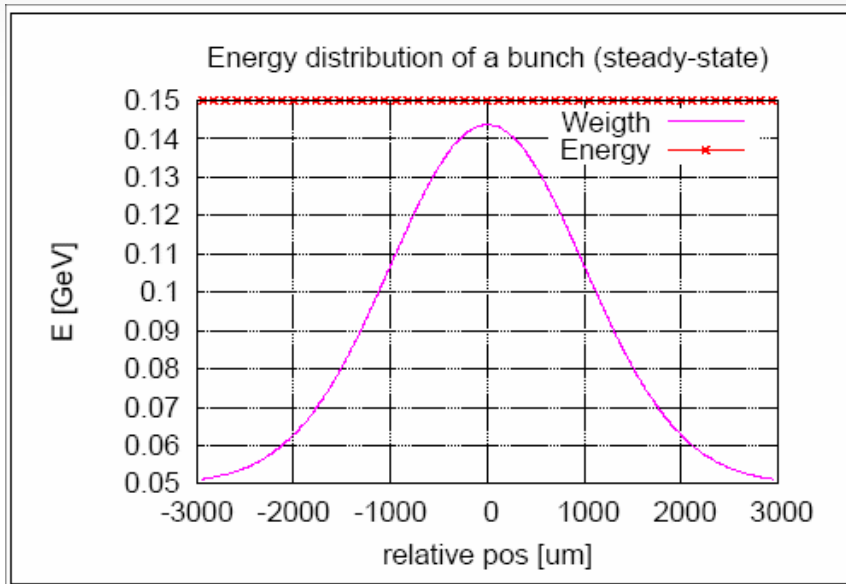
- For our initial parameters we get

$$r_{ad,after} = 8.3mm, r_{ad,initial} = 5.0mm$$

- Meaning: with the nominal parameters cited above we will have a resulting 3σ beam size of 8.3mm due to the adiabatic undamping alone (while half-aperture is $a_0=11.5mm$) !

Beam envelope along the lattice

- Thus, beam envelope along the lattice $r_{ad} \propto 1/\sqrt{\gamma}$,
 γ for lowest particle



The CTF3 drive beam

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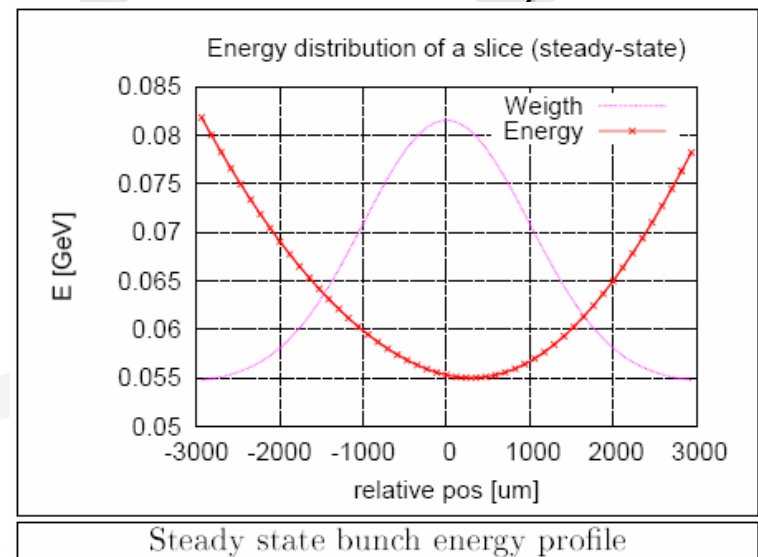
- Resulting parameters:
 - $P = 166$ MW (steady-state power production)
 - $S = 63.3$ % (max .energy spread)
 - $\eta = 61.3$ % (steady-state extraction efficiency)
 - $r_{ad} = 8.3$ mm (**3-sigma envelope due to adiabatic effects alone**)

Smaller beam envelope?: reduce the current

- The adiabatic envelope, r_{ad} , can be made as small (large) as we want by decreasing (increasing) the current, I .
- For the TBL, I can be calculated on the fly as function of r , E_0 , N and PETS parameters (only because: $\eta \approx S \times F(\sigma)$, $\eta_{dist} \approx 1$)
- (Calcs omitted)
- However, decreasing the current will mean less power extracted, P , and less extraction efficiency η achieved (while we want to show as high P and η as possible)
- E.g. for if we want a $r_{ad} = (2/3)a_0 = 7.7\text{mm}$ we must reduce current to $I=27\text{A}$ (with a corresponding lower $P=135\text{ W}$, $\eta = 55\%$)

Challenge: beam dynamics calculations

- We have up to 90% energy spread S (CLIC)
- Spread acts stabilizing (different betatron wavelength lead to decoherence of transverse kicks) – but difficult to calculate the effect
- Also the finite group velocities and damping makes calculations difficult
- No analytical formulas or framework available (ongoing work, try to get somewhere, but no results so far)
- → **need for simulations**

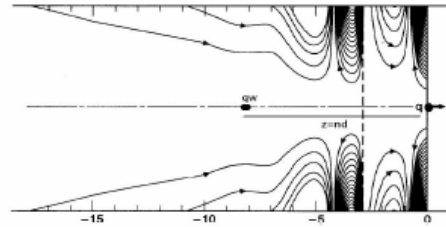


Part 2

Effect of transverse wakes (summary)

Transverse wakes

- A source particle q_s induces wake fields in PETS cavity



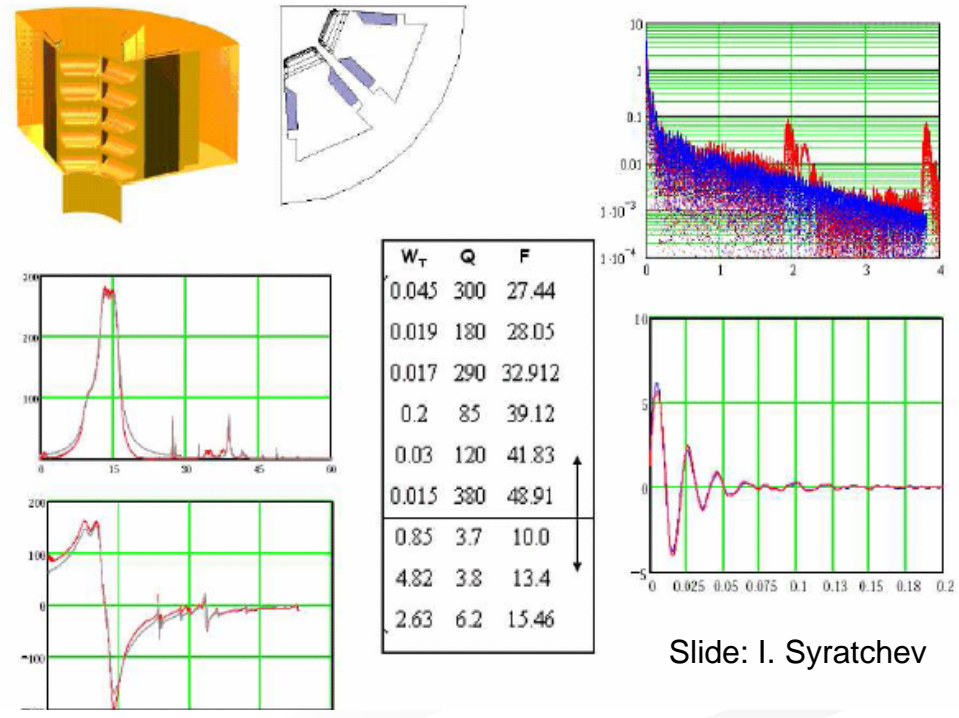
- A witness particle q_w , following at a distance z , is kicked by the fields from leading particles
- The total transverse force on q_w is given by (1D)

$$\int_0^{l_{cav}} F_y(z) ds \approx -\Delta y q_s q_w w_T(z)$$

where $w_T(z)$ is the transverse dipole wake function - the “ δ -wake” (h.o.ms ignored here)

PLACET input: dipole wake function

- PETS are modelled with GdfidL (I. Syratchev)
- For a given PETS structure, the transverse δ -wake / impedance is calculated



PLACET simulations

- Multiple modes identified from GdfidL calc
- For each mode, $w_{T_i}, Q_i, f_{T_i}, \beta_{T_i}$ are identified
- The total wake function for each mode thus:

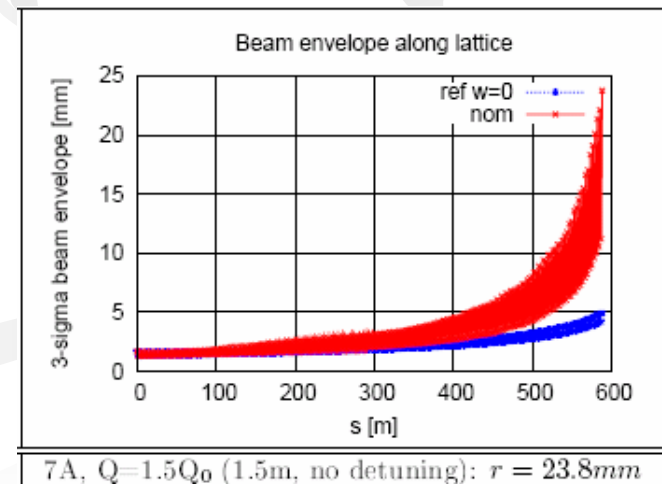
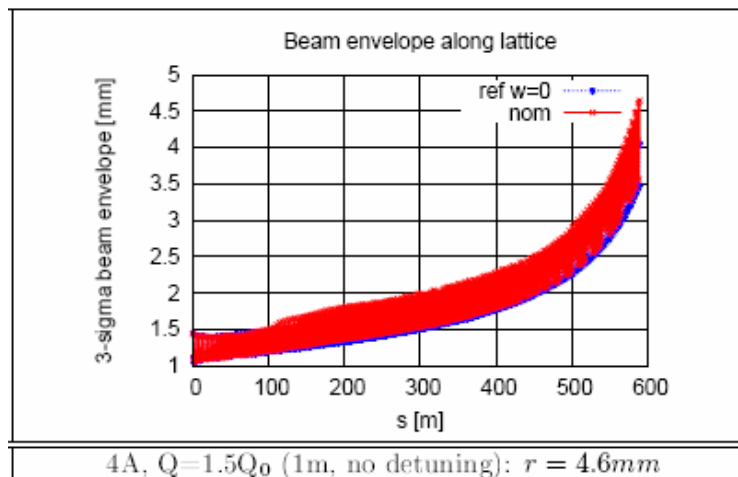
$$W_{T_i}(z) = w_{T_i} \sin\left(\omega \frac{z}{c}\right) \left(L - z_{ij} \frac{\beta_T}{1 - \beta_T}\right) e^{-z\omega/2cQ(1-\beta_T)} [V/Cm]$$

- Transverse kick of q_w :

$$\Delta y'_w = \sum_{modes} \frac{\Delta p_{y,w}}{m_w c} = \sum_{modes} y_s \frac{q_s q_w}{E_w} W_{\delta T}(z) [rad]$$

Goal: transverse wakes should not amplify beam jitter

- A design target for the PETS is to ensure that beam jitter are not amplified significantly due to transverse wakes (and leading to beam blow-up)
- A number of simulations has been run (initiated by I. Syratchev)
- Results: basically no problem for *nominal* PETS parameters (both CLIC and TBL lattice checked)
- (Example: CLIC low β /high β FODO lattice)

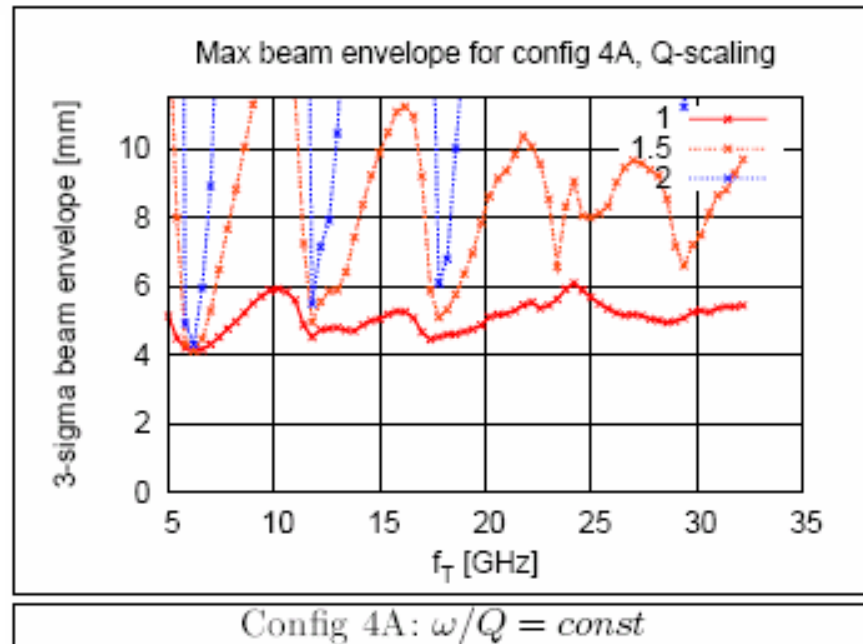


More examples of PETS test simulations

- Metric used: 3 – *sigma* beam envelope at end of lattice
- Initial conditions: beam with initial static offset + jitter at the transverse resonance frequency
- Beam blow-up depends on z/λ_{T_i} : $\sin(\frac{2\pi}{\lambda_{T_i}}z) = 0 \Rightarrow z = \frac{n}{2}\lambda_T \Rightarrow f_T = \frac{n}{2}12GHz$ (zeros)

SUM mode, $\Sigma w, \bar{Q}$

$w = 8.3, Q = 4.6$

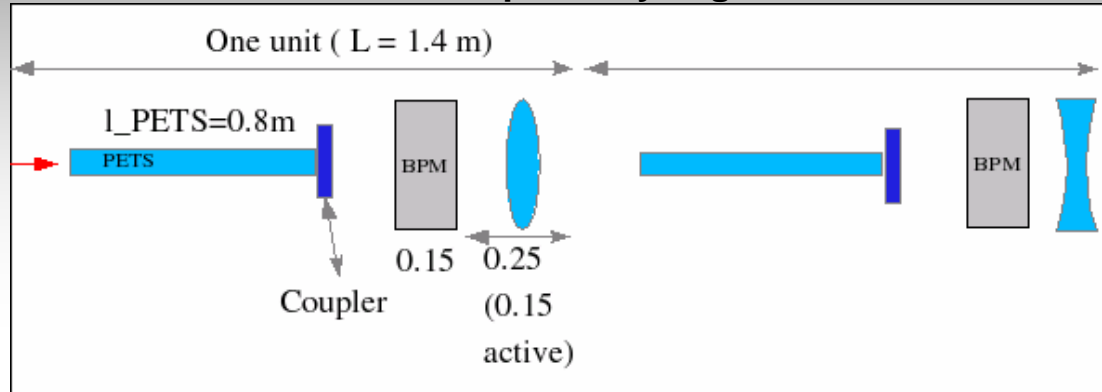


Part 3

Effect of component misalignment

5 types of misalignment studied

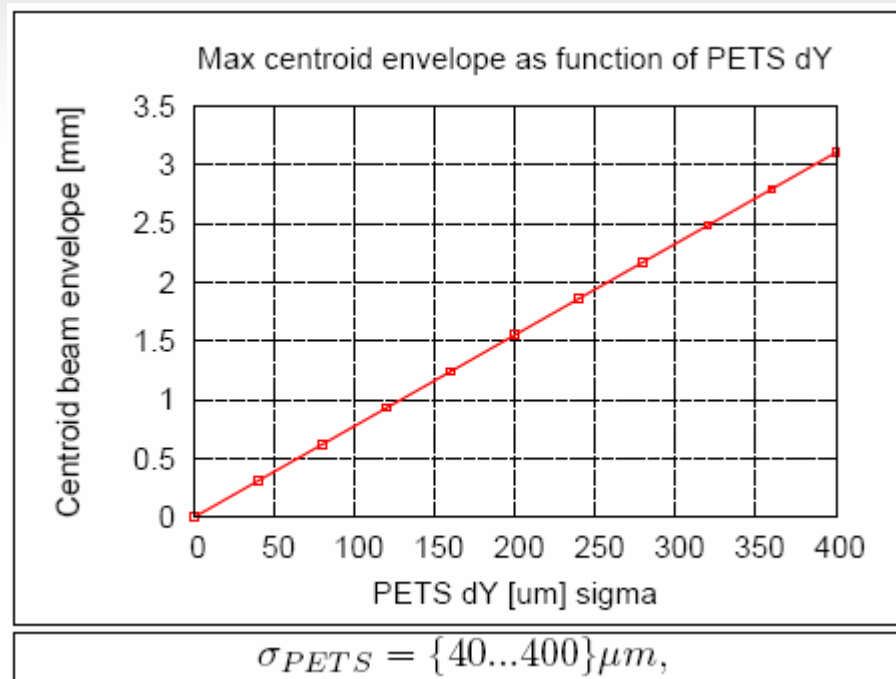
- In part 1 and 2: all simulations was with a **perfectly aligned machine**



- Now: will study **the effect of misalignment** of machine components
- Each misalignment (PETS, Quads) studied separately
- 100 random machines simulated for each case. Metric: max. **centroid** offset, r_c , along lattice (of all machines)
- The initial beam will be assumed to be on the reference orbit
- Adiabatic effect is **NOT** included in order to study each effect separately (no macroparticles distribution)
- Total 3-sigma beam envelope will therefore be
 - 3-sigma adiabatic envelope "+" centroid envelope : r_{ad} "+" r_c (where the "+" is only in worst case a real +)
- Still: with the adiabatic envelope $r = 8.3$ mm (versus half-aperture of $a_0 = 11.5$ mm) **we do not have a large "envelope budget" for component misalignment**

1) Position offset of PETS

- A PETS off axis will induce transverse kicks (dipole wake $\propto y_{\text{source}}$)
- We scatter the PETS in y: $\sigma_{\text{PETS}} = \{40 \dots 400\} \mu\text{m}$

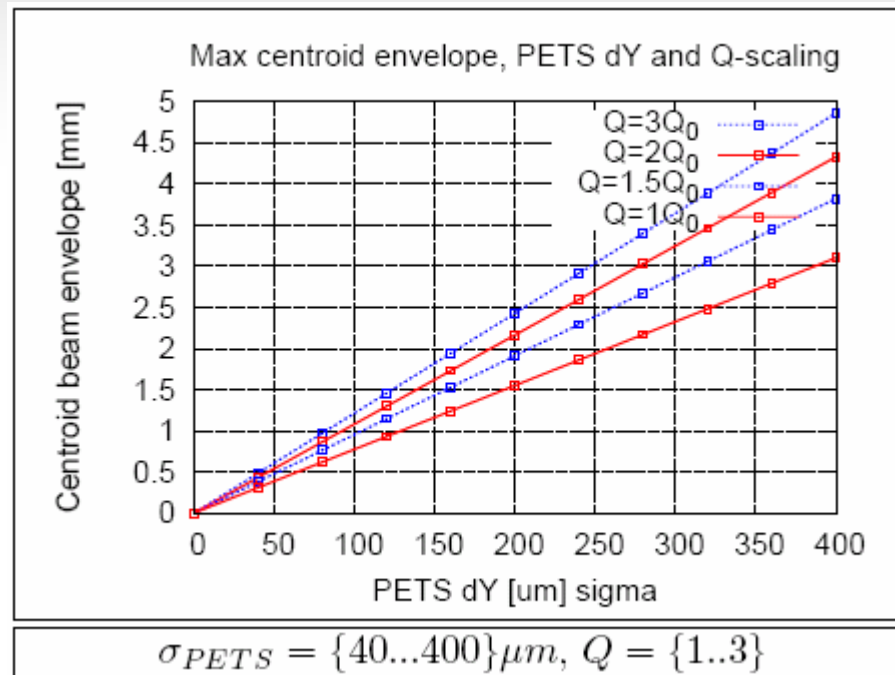


(linear graph due to the linear lattice model and same seeds in all simulations - all info in one point)

- Prelim. criterion: centroid envelope $< 1 \text{ mm}$
 $\Rightarrow \sigma_{\text{PETS}} < 120 \mu\text{m}$

Position offset of PETS with Q-scaling

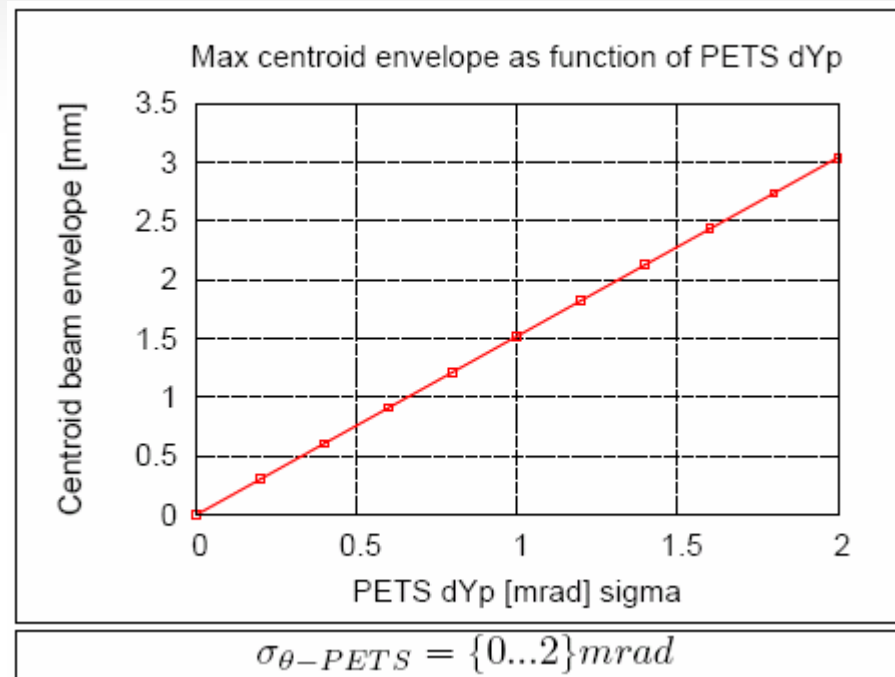
- Also interesting to see effect of large Q in this scenario (previous PETS simulations imply that the effect should not be drastic)



- We see that as long as $\sigma_{PETS} < 100 \mu\text{m}$ we are still OK, even for a factor $Q=3Q_0$
- (But this is not “worst case scenario” for PETS: jitter on resonance)

2) Angle offset of PETS

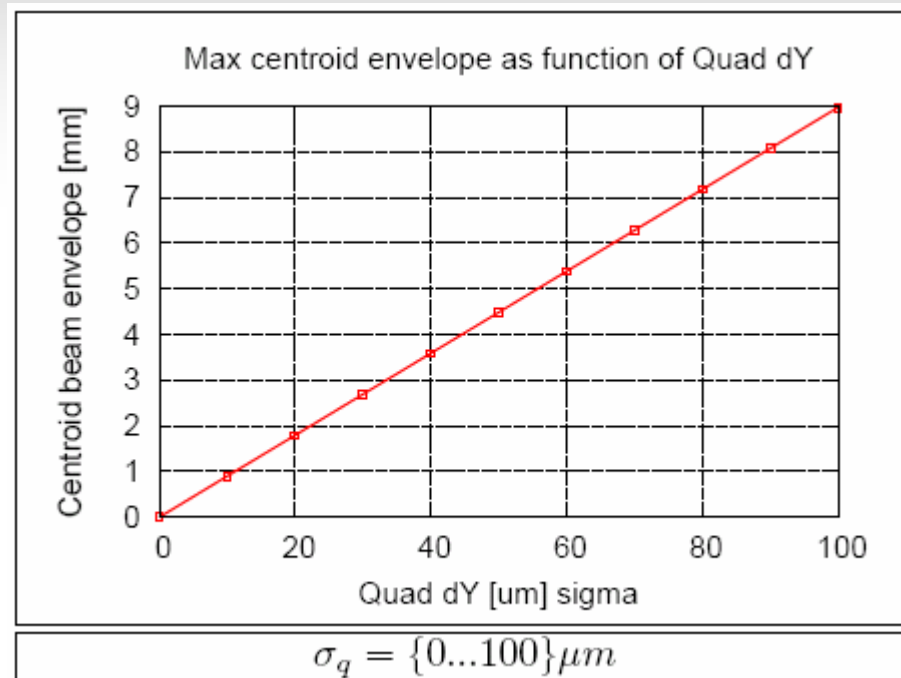
- An angle offset (around x,y) of the PETS centre should basically have the same effect as the corresponding position offset, $\sigma_{\theta-PETS} = (\sigma_{PETS} / 0.5l_{PETS}) * 2$. Just to confirm:



- Prelim. criterion: centroid envelope < 1 mm
 $\Rightarrow \sigma_{\theta-PETS} < 0.6 \text{ mrad}$
- (An angle offset around s: negligible effect)

3) Position offset of quadrupoles

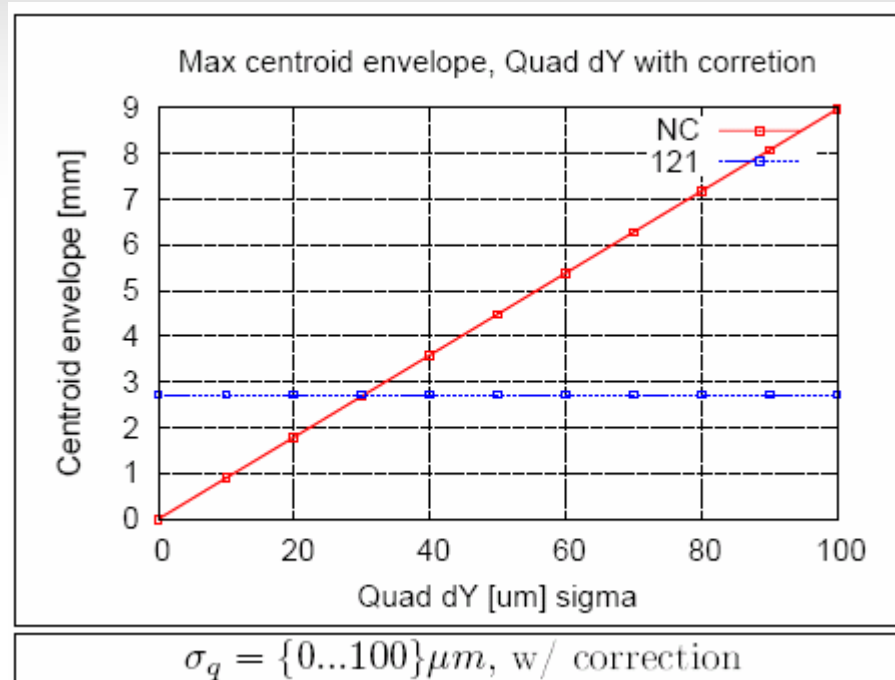
- An offset in a quad will induce a dipole kick
- We scatter the quads in y: $\sigma_q = \{10 \dots 100\} \mu\text{m}$



- We see that without any correction of the quad positions we should require a final alignment $\sigma_{\text{PETS}} < 10 \mu\text{m}$ (probably not feasible with prealignment ?)
- Solution: beam-based alignment (BBA)

BBA for quads: 1-to-1 correction

- Quads in TBL (as in CLIC) will be on movers
- The simplest BBA: steer each quad so that the beam goes through centre of the following BPM



- Implementation: correction can in principle be performed in one go. **b**: BPMs after one pulse. **R**: the response matrix: $\Delta \mathbf{y}_c = \mathbf{R} \mathbf{b}$
- Effectively: **quad position error σ_q , is transferred to BPM position error σ_{BPM}**
- 1-to-1 correction can be needed as a first correction – but we can do better

BBA for quads: dispersion free steering

- The dipole kicks resulting from quad offset will induce dispersion (in the sense “energy-dependent trajectories”) in the lattice
- Idea: move quads so that beams of different initial energies follows the same trajectory
- E.g. send the nominal beam with E_0 and a test-beam with $E_1=0.8\times E_0$
- This can, in principle, be implemented by generating the response matrix of both the nominal beam, \mathbf{R}_1 , and for a beam of e.g. lower energy, \mathbf{R}_2 , measure and do the correction $\Delta\mathbf{y}_c=(\mathbf{R}_1 - \mathbf{R}_2)^+(\mathbf{b}_1-\mathbf{b}_2)$ (weighted against 1-to-1 correction)
- Effectively: **quad position error σ_q , is transferred to BPM resolution error σ_{res}**

Special variant of DFS needed for TBL

- TBL and CLIC deceleration station:
 - cannot use lower energy beam due to beam stability
 - higher initial energy beam not available
- Trick: we can use a beam with lower current instead
Wakefields will be lower and beam will quickly have higher energy
- Can either reduce bunch charge, or take out a number of bunches (probably easier?)

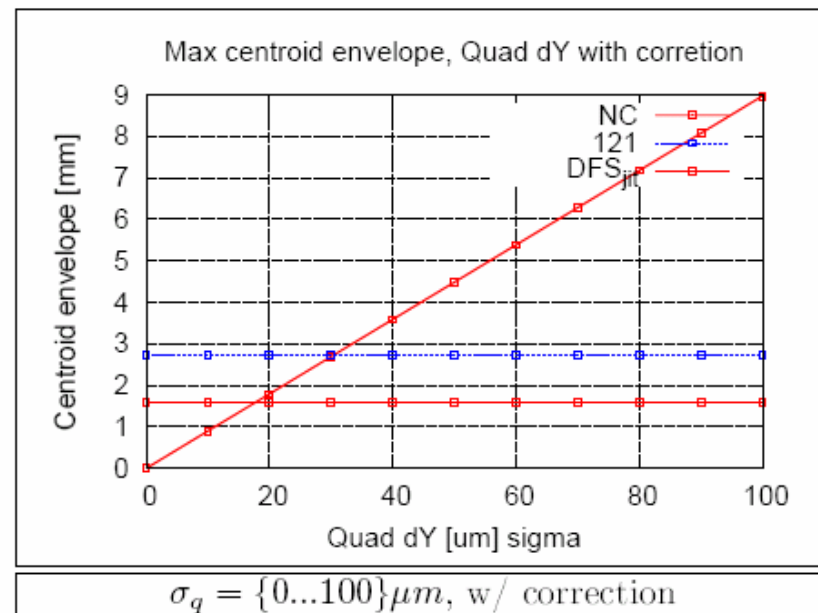
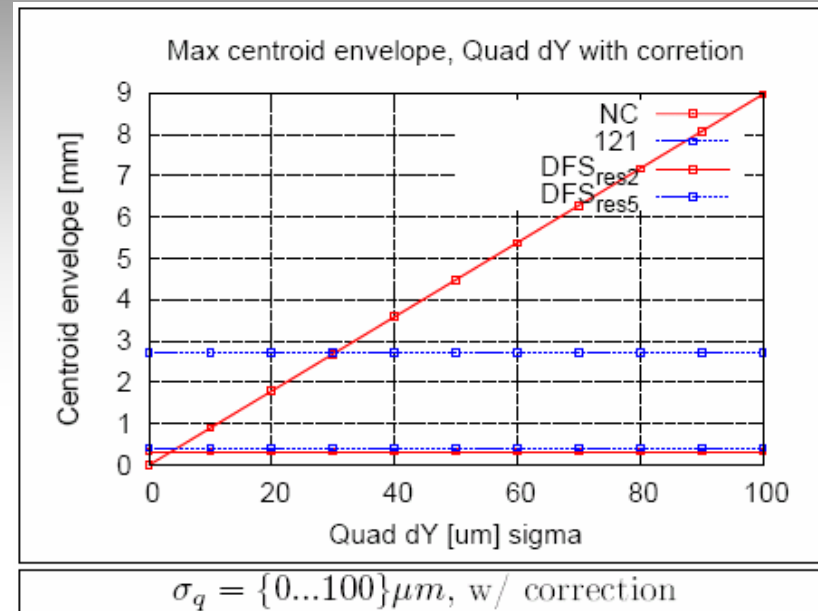


- Results seems to be as good as for energy test beam
- All results from now are run with optimal DFS weighting

NB: now very easy to test new BBA algorithms in PLACET with then new Octave interface (A. Latina)

Correction with current test beam DFS

- Seemingly very good result (with our linear lattice and ideal elements), and shows the principle
- Remains to be studied: the DFS algorithm is, among other things, sensitive to jitter in the main/test beams. E.g. inducing uncorrelated random jitter $\sigma_{\text{beam}} = 100\mu\text{m}$ on both beams gives a substantially worse result. However, this can be partially remedied (D. Schulte), but not studied further here.
- In the worst case: 1-to-1 correction would give as good result as DFS

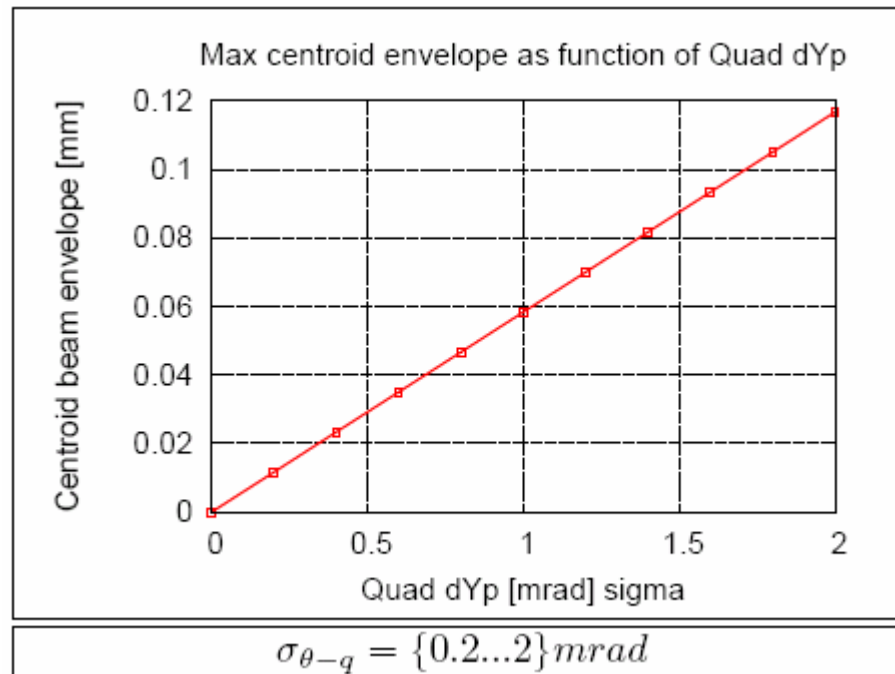


TBL as a test-bed for the CLIC station

- In any case: it would be very interesting to use the TBL to test the concept of dispersion free steering and other beam-based alignment
- One would like to have similar conditions to the CLIC station (ideally: similar alignment precision and BPM resolution)
- TBL: also test-bed for automation of BBA ?

4) Angle offset of quadrupole

- A small rotation around s (skew quadrupole) will have negligible effects for the 8 FODO cell lattice in question
- A small rotation around x, y will have negligible effect (integrated force ≈ 0 , due to negligible motion in a quad)
- Quad angle offset in y: $\sigma_{\theta q} = \{0 \dots 2\}$ mrad . Just to confirm:



- Negligible for even relatively large rotation

Discussion

- Trustworthiness of the simulation results

“Are we now sure to get the TBL beam through???”

- In “favour”
 - Our metrics are conservative
- In “disfavour”
 - several important idealizations:
 - lattice: except scattering and BPM res: ideal lattice elements
 - lattice: linear lattice
 - beam: Initial 0 bunch energy spread
 - beam: gaussian bunch shape
 - wakes: Only monopole and dipole wake
 - wakes: PETS Q factor difficult to predict
 - ...and more
- First conclusion: a TBL is still needed...

Conclusions

- The nominal PETS parameters given seem satisfactory
- We see that the difference between a position accuracy of $\sigma = 100 \mu\text{m}$ and $\sigma = 200 \mu\text{m}$ is significant
- Under the conditions in these simulation, with $\sigma = 200 \mu\text{m}$ we cannot get the beam fully through even when doing 1-to-1 correction (only with perfect DFS steering) (without reducing I, P and η)
- The margin for $\sigma = 100 \mu\text{m}$ is much more comforting. Ideally :
 - ⇒ $\sigma_{\text{PETS}} \leq 100 \mu\text{m}$ (minimize transverse wake kick)
 - ⇒ $\sigma_{\theta\text{-PETS}} \leq 0.5 \text{ mrad}$ (same reason)
 - ⇒ $\sigma_{\text{quads}} \leq 100 \mu\text{m}$ (get beam through in order to prepare for further corrections)
 - ⇒ $\sigma_{\text{BPM}} \leq 100 \mu\text{m}$ and $\sigma_{\text{res}} \leq 5 \mu\text{m}$
 - ⇒ Mover positioning resolution in x and y: same O.M. as σ_{res}